



Recent results on francium isotopes at the collinear resonance ionisation spectroscopy (CRIS) setup.

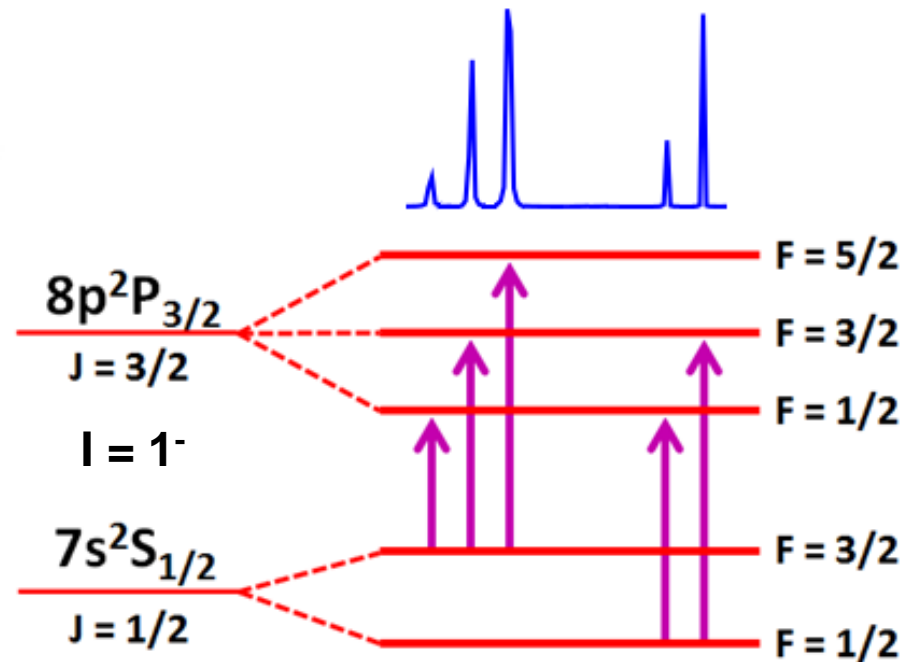
Gregory Farooq-Smith

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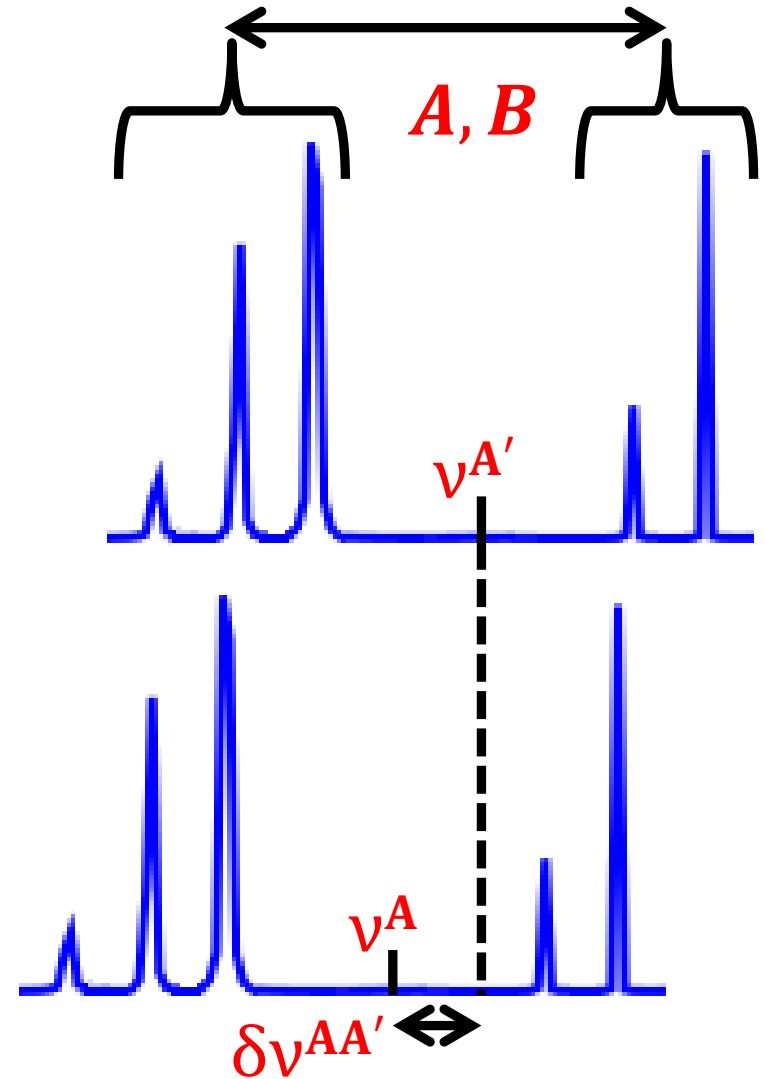
Laser Spectroscopy

- Hyperfine spectra can be examined by scanning the atomic transitions of an isotope with a laser.
- The extracted parameters can be related to fundamental nuclear properties of the isotope.

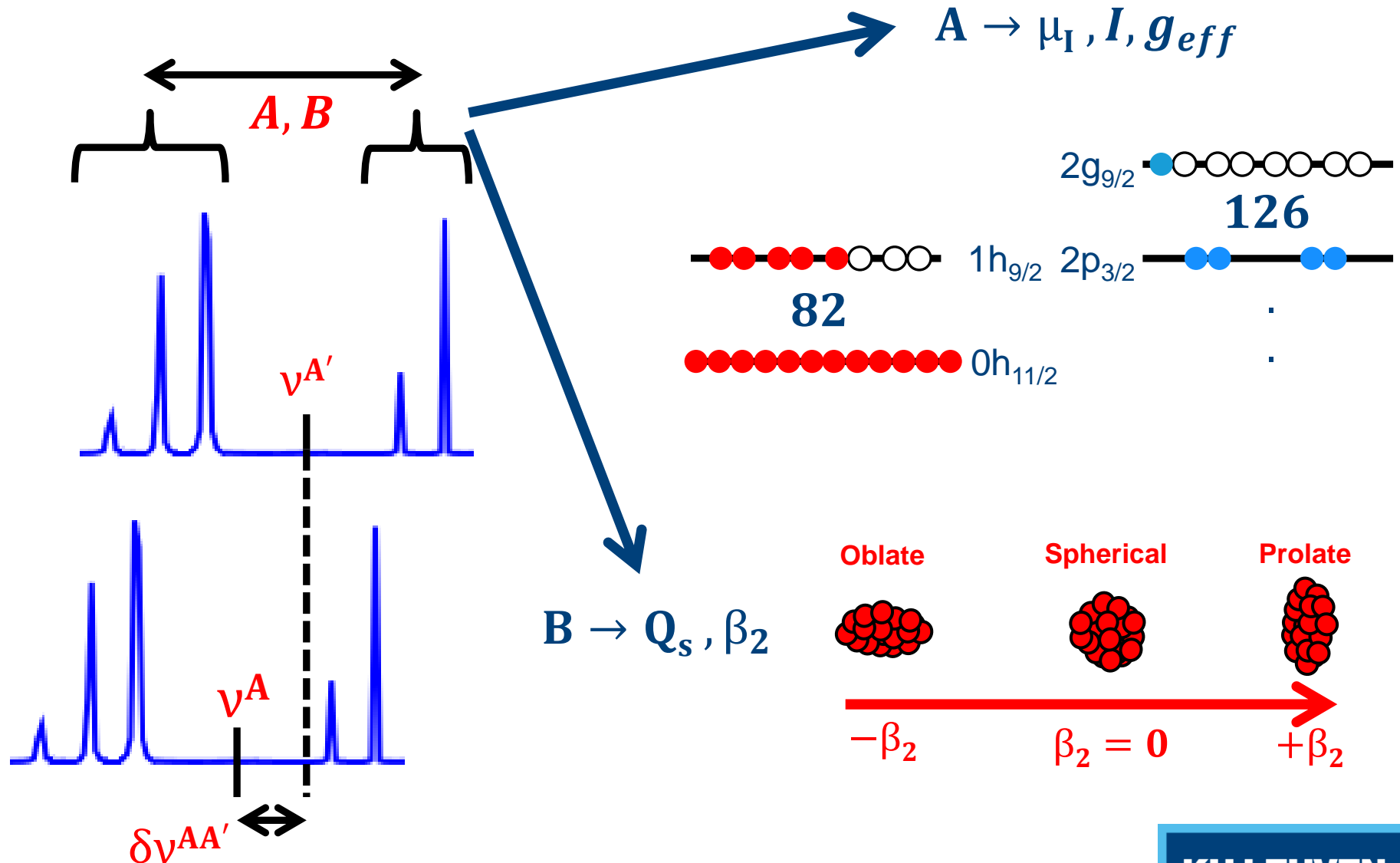


Laser Spectroscopy

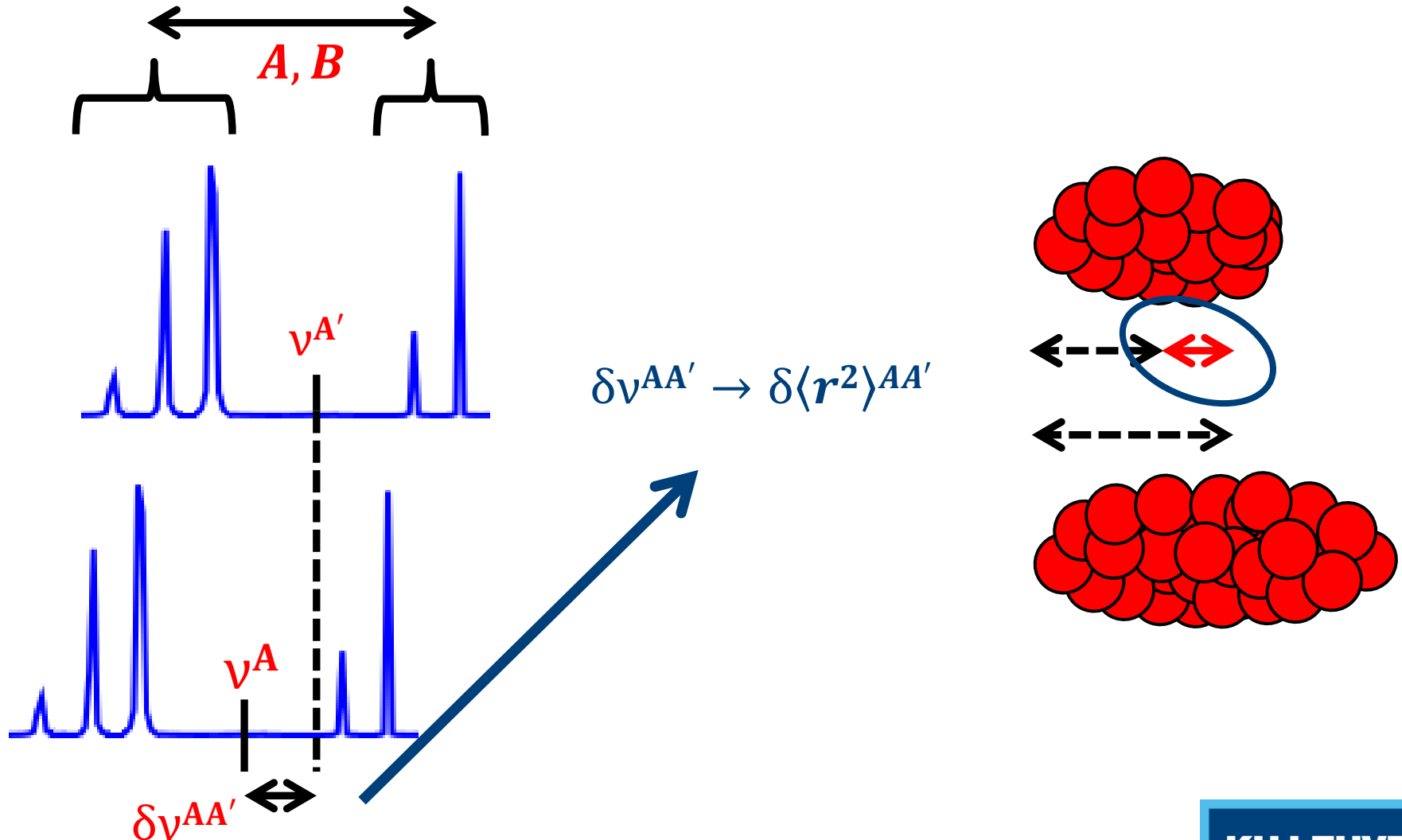
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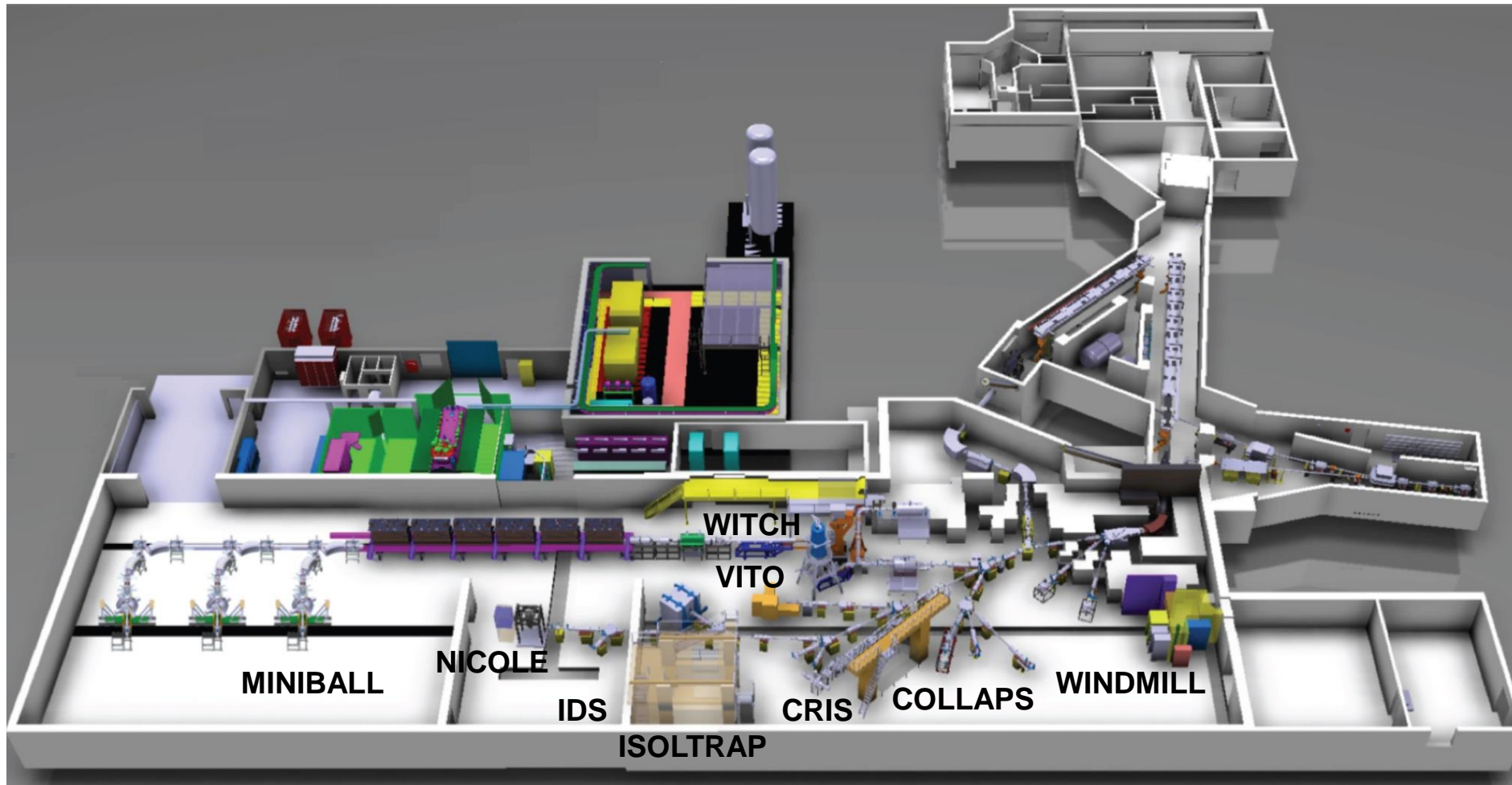
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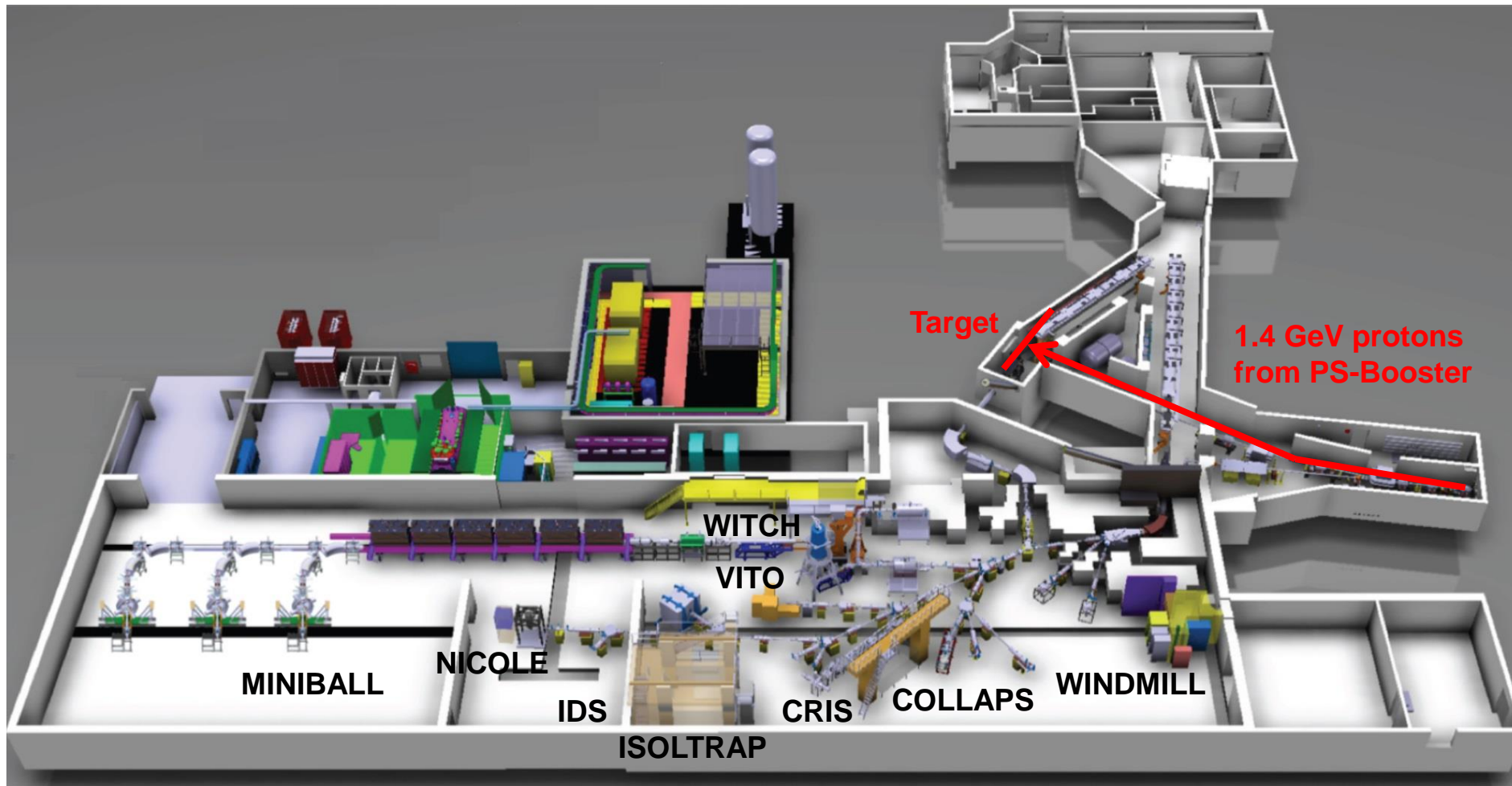
Laser Spectroscopy



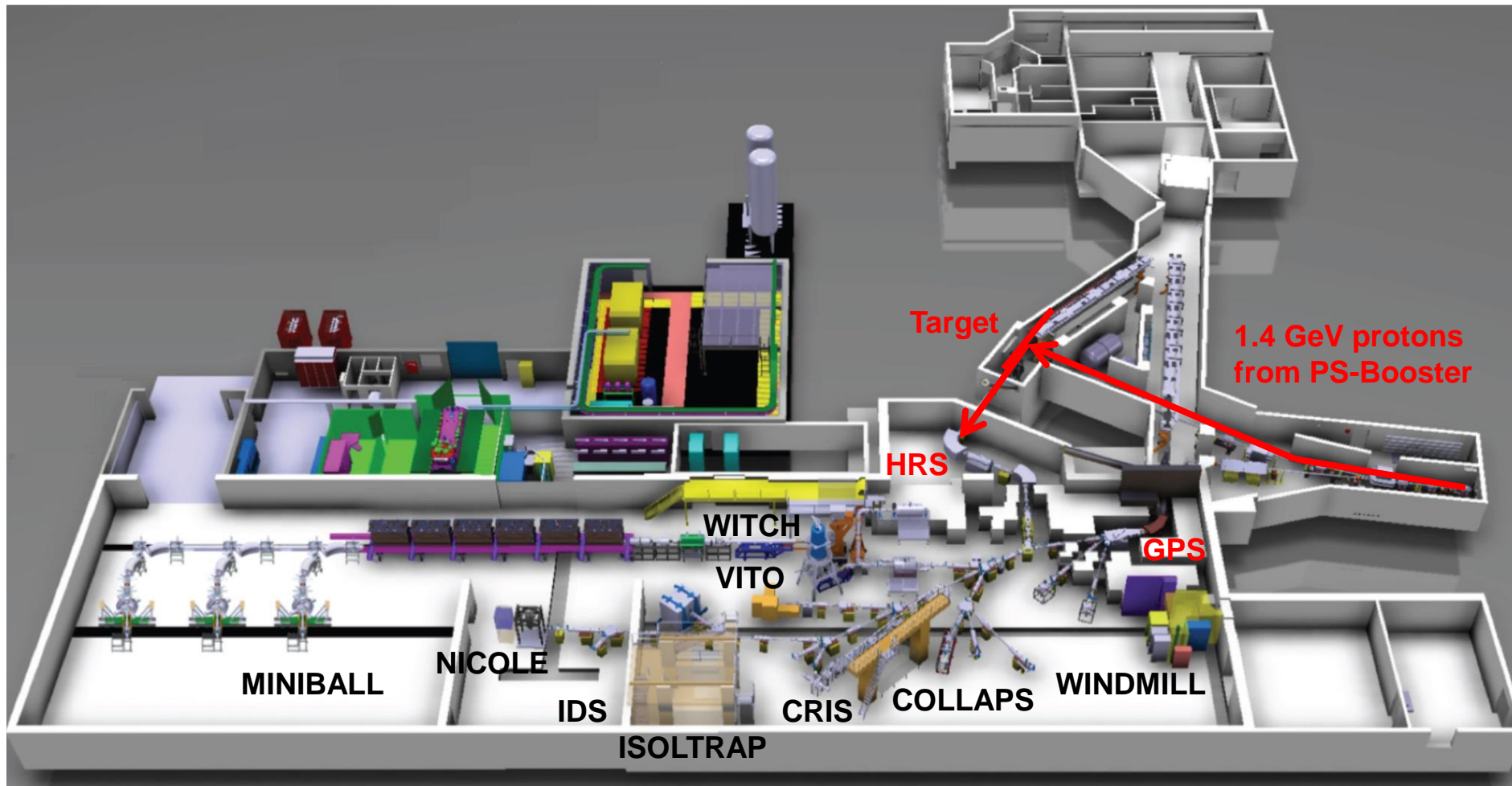
CRIS, CERN-ISOLDE



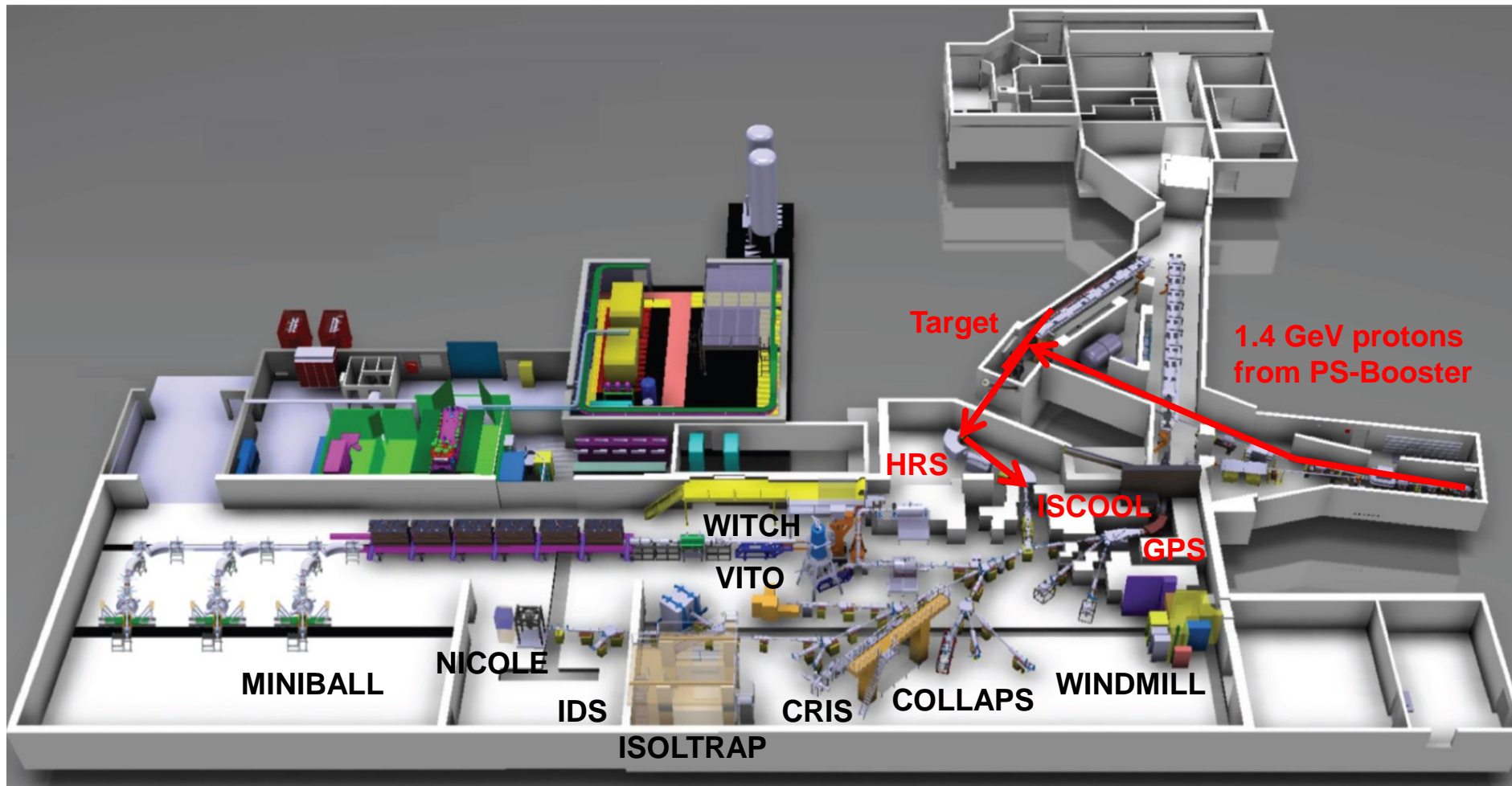
CRIS, CERN-ISOLDE



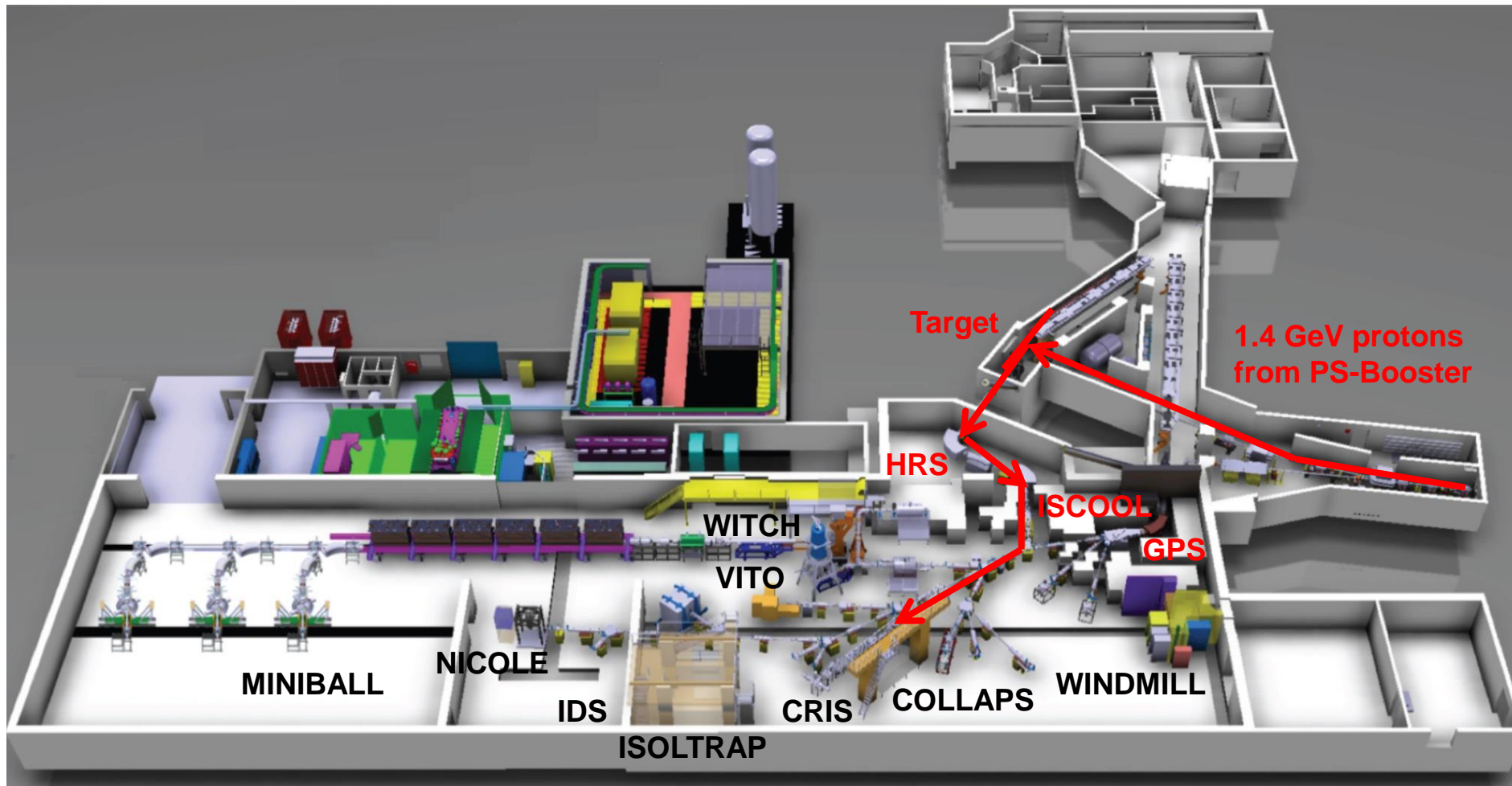
CRIS, CERN-ISOLDE



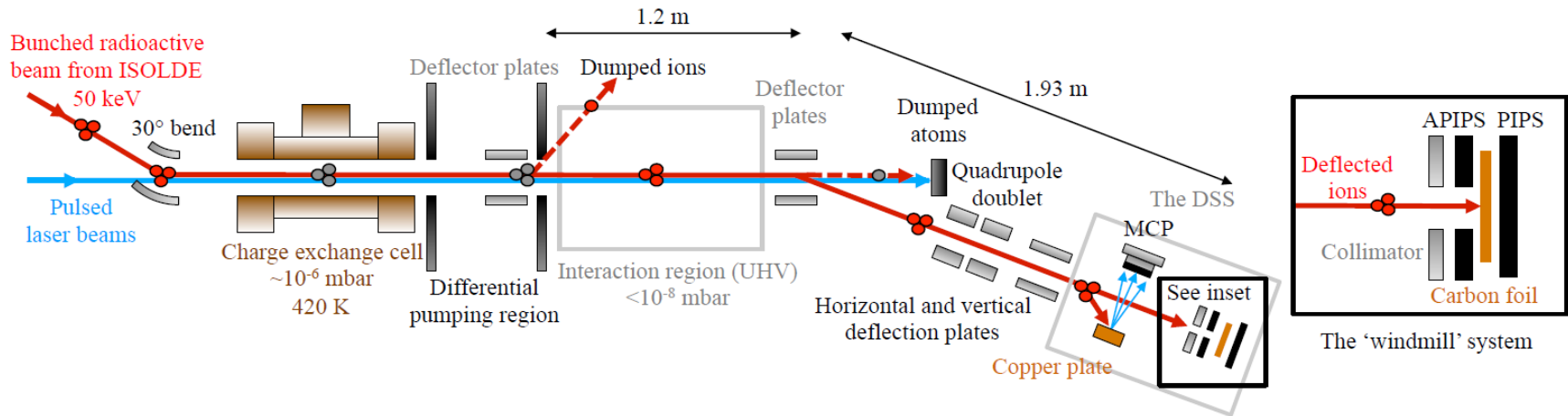
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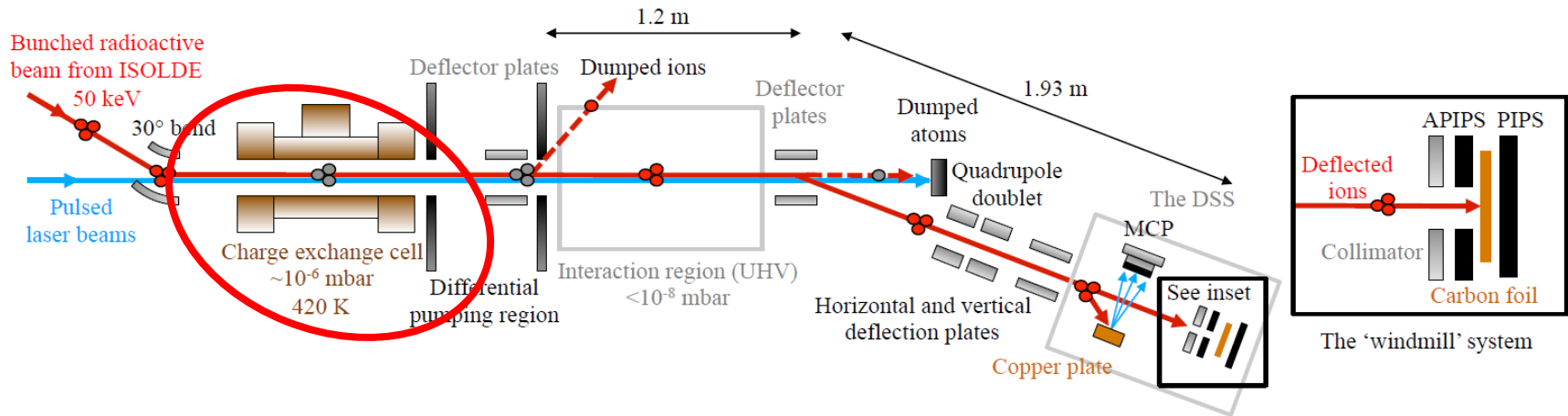
CRIS, CERN-ISOLDE



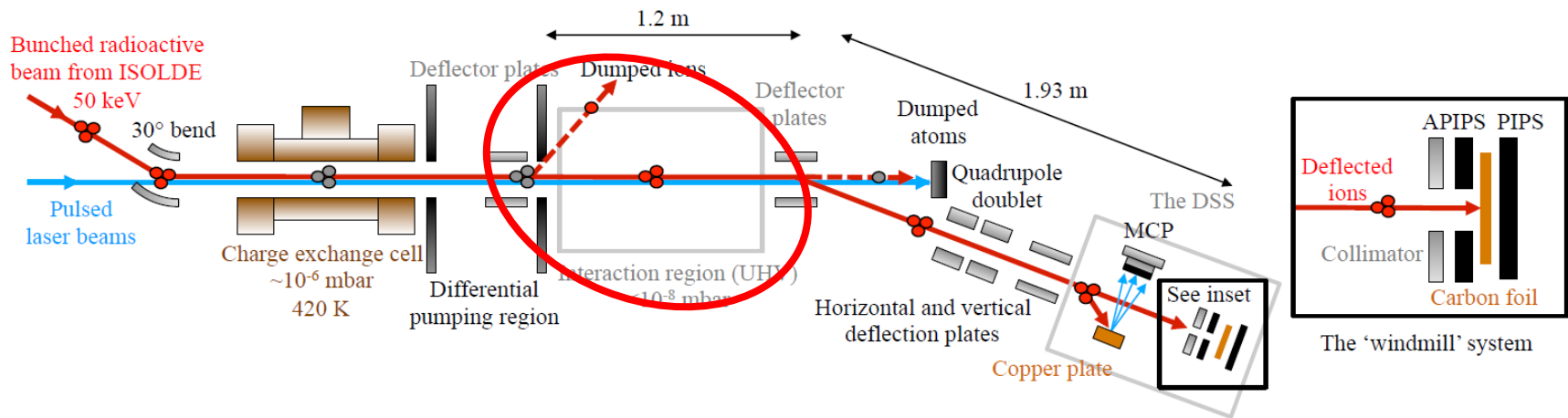
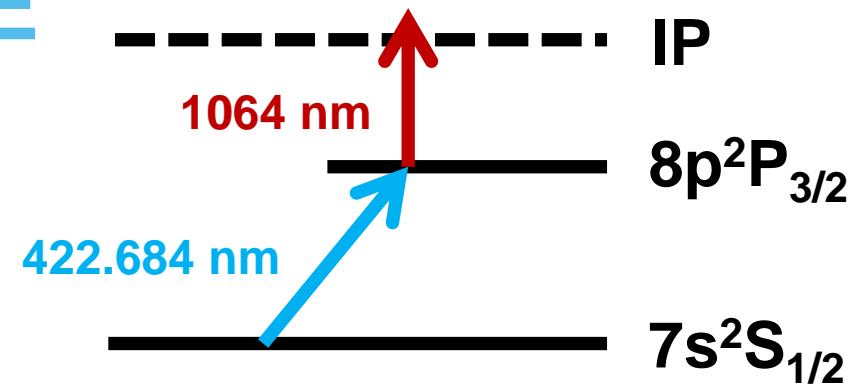
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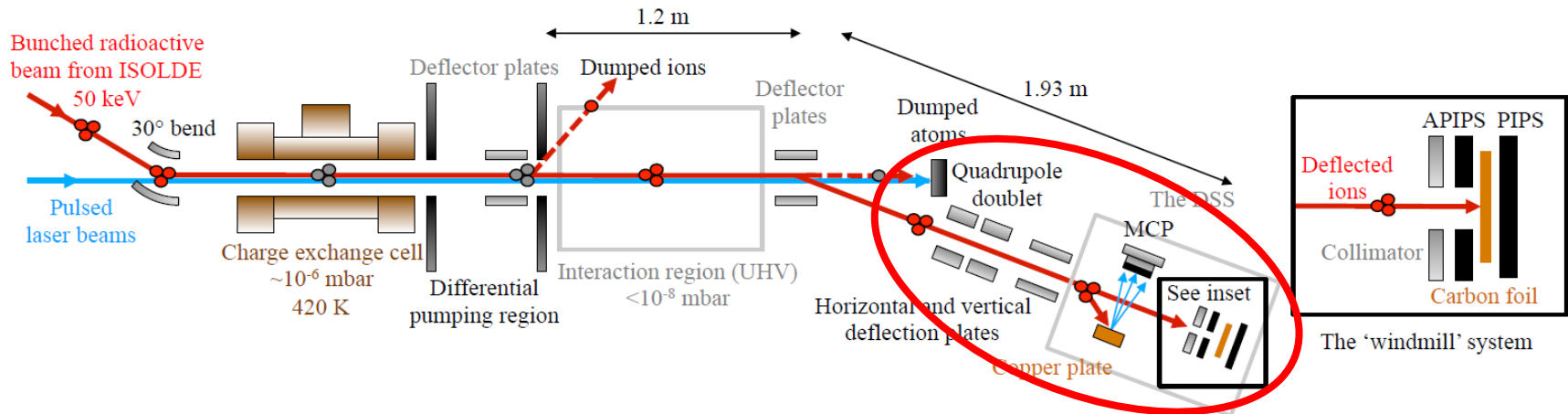
CRIS, CERN-ISOLDE



CRIS, CERN-ISOLDE



CRIS, CERN-ISOLDE



IS471: Why francium?

- An ongoing experimental campaign is investigating francium isotopes at the CRIS setup.
- It is thought that a variety of nuclear features exist in this isotopic region, depending on the proximity to $Z = 82$.

	201Ra	202Ra	203Ra	204Ra	205Ra	206Ra	207Ra	208Ra	209Ra	210Ra	211Ra	212Ra	213Ra	214Ra	215Ra	216Ra	217Ra	218Ra	219Ra	220Ra	221Ra	222Ra	223Ra	224Ra	225Ra	226Ra	227Ra	228Ra	229Ra	230Ra	231Ra	232Ra	233Ra	234Ra
199Fr	200Fr	201Fr	202Fr	203Fr	204Fr	205Fr	206Fr	207Fr	208Fr	209Fr	210Fr	211Fr	212Fr	213Fr	214Fr	215Fr	216Fr	217Fr	218Fr	219Fr	220Fr	221Fr	222Fr	223Fr	224Fr	225Fr	226Fr	227Fr	228Fr	229Fr	230Fr	231Fr	232Fr	233Fr
196Rn	199Rn	200Rn	201Rn	202Rn	203Rn	204Rn	205Rn	206Rn	207Rn	208Rn	209Rn	210Rn	211Rn	212Rn	213Rn	214Rn	215Rn	216Rn	217Rn	218Rn	219Rn	220Rn	221Rn	222Rn	223Rn	224Rn	225Rn	226Rn	227Rn	228Rn	229Rn	230Rn	231Rn	
197Ac	198Ac	199Ac	200Ac	201Ac	202Ac	203Ac	204Ac	205Ac	206Ac	207Ac	208Ac	209Ac	210Ac	211Ac	212Ac	213Ac	214Ac	215Ac	216Ac	217Ac	218Ac	219Ac	220Ac	221Ac	222Ac	223Ac	224Ac	225Ac	226Ac	227Ac	228Ac	229Ac		
196Po	197Po	198Po	199Po	200Po	201Po	202Po	203Po	204Po	205Po	206Po	207Po	208Po	209Po	210Po	211Po	212Po	213Po	214Po	215Po	216Po	217Po	218Po	219Po	220Po	221Po	222Po	223Po	224Po	225Po	226Po	227Po			
195Bi	196Bi	197Bi	198Bi	199Bi	200Bi	201Bi	202Bi	203Bi	204Bi	205Bi	206Bi	207Bi	208Bi	209Bi	210Bi	211Bi	212Bi	213Bi	214Bi	215Bi	216Bi	217Bi	218Bi	219Bi	220Bi	221Bi	222Bi	223Bi	224Bi					
194Pb	195Pb	196Pb	197Pb	198Pb	199Pb	200Pb	201Pb	202Pb	203Pb	204Pb	205Pb	206Pb	207Pb	208Pb	209Pb	210Pb	211Pb	212Pb	213Pb	214Pb	215Pb	216Pb	217Pb	218Pb	219Pb	220Pb								
193Tl	194Tl	195Tl	196Tl	197Tl	198Tl	199Tl	200Tl	201Tl	202Tl	203Tl	204Tl	205Tl	206Tl	207Tl	208Tl	209Tl	210Tl	211Tl	212Tl	213Tl	214Tl	215Tl	216Tl	217Tl										
192Hg	193Hg	194Hg	195Hg	196Hg	197Hg	198Hg	199Hg	200Hg	201Hg	202Hg	203Hg	204Hg	205Hg	206Hg	207Hg	208Hg	209Hg	210Hg	211Hg	212Hg	213Hg	214Hg	215Hg	216Hg										

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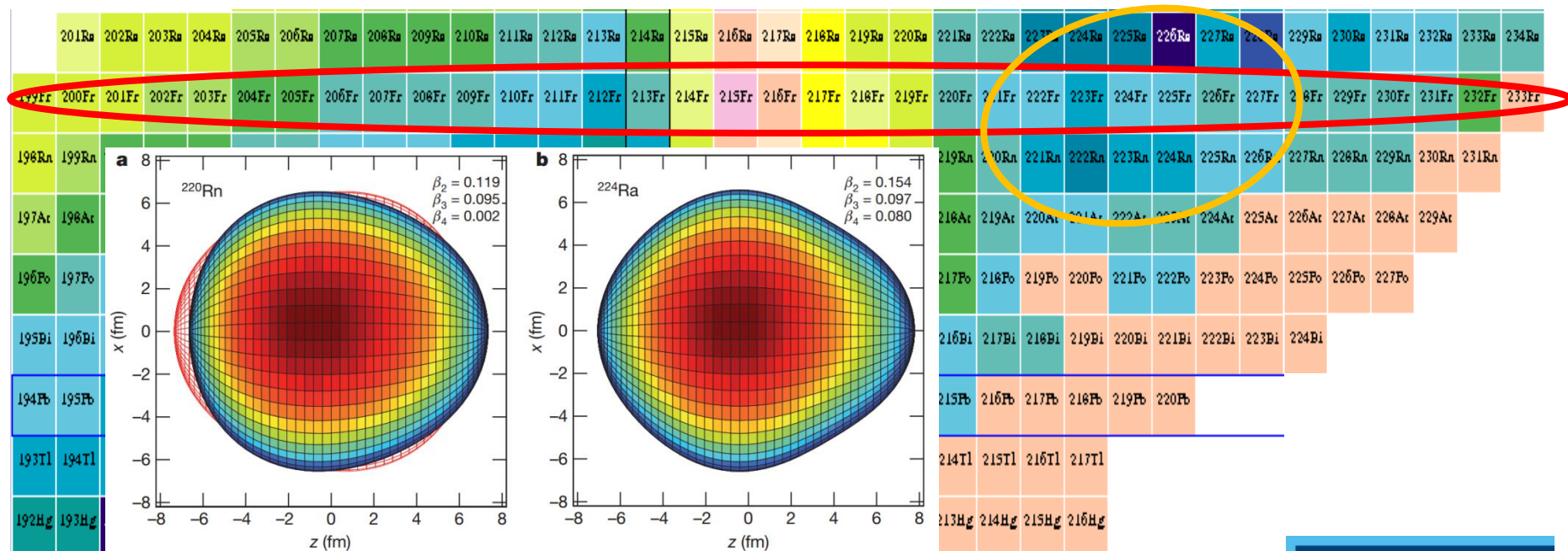
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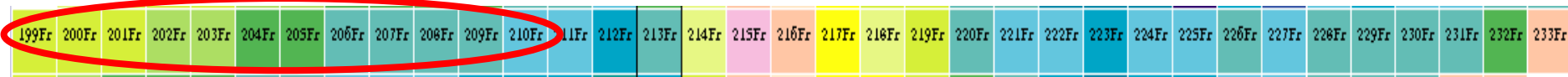
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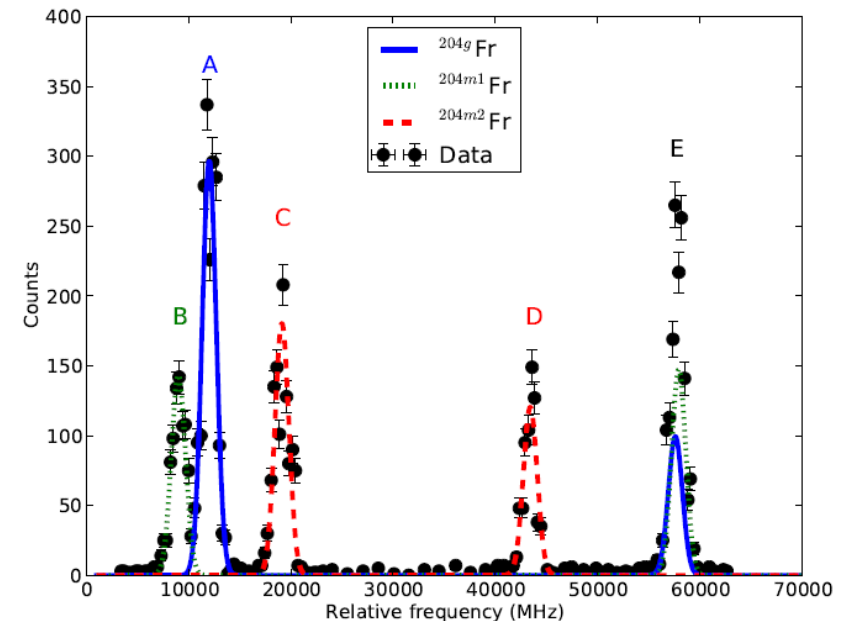
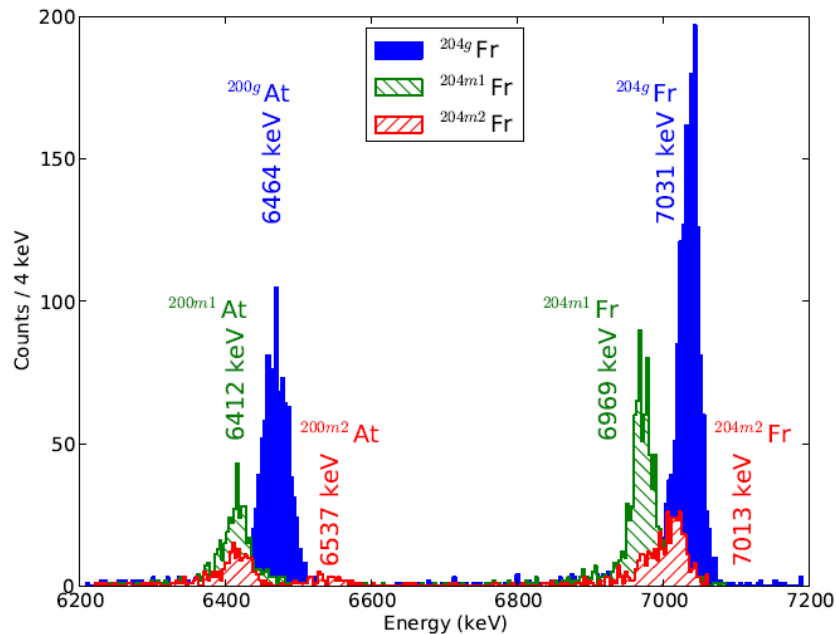
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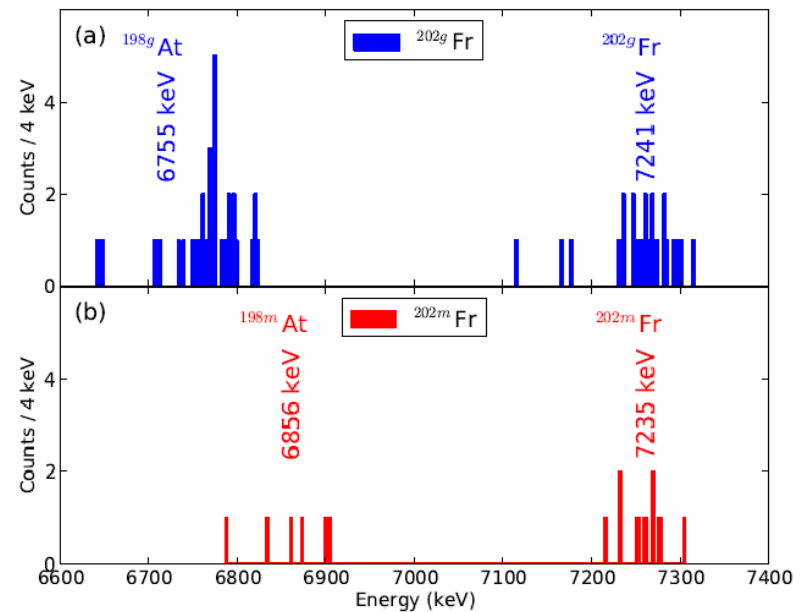
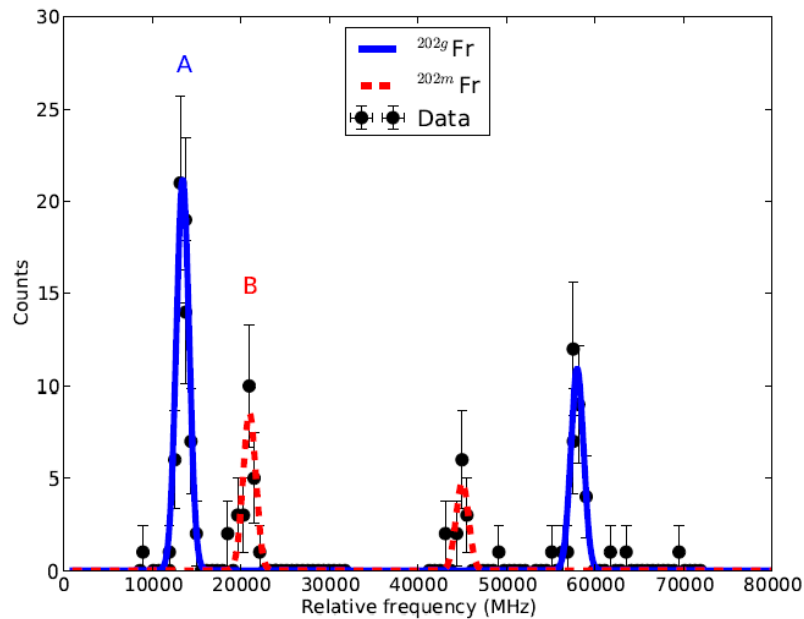
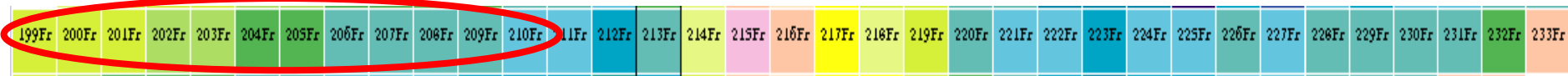
IS471: Neutron deficient



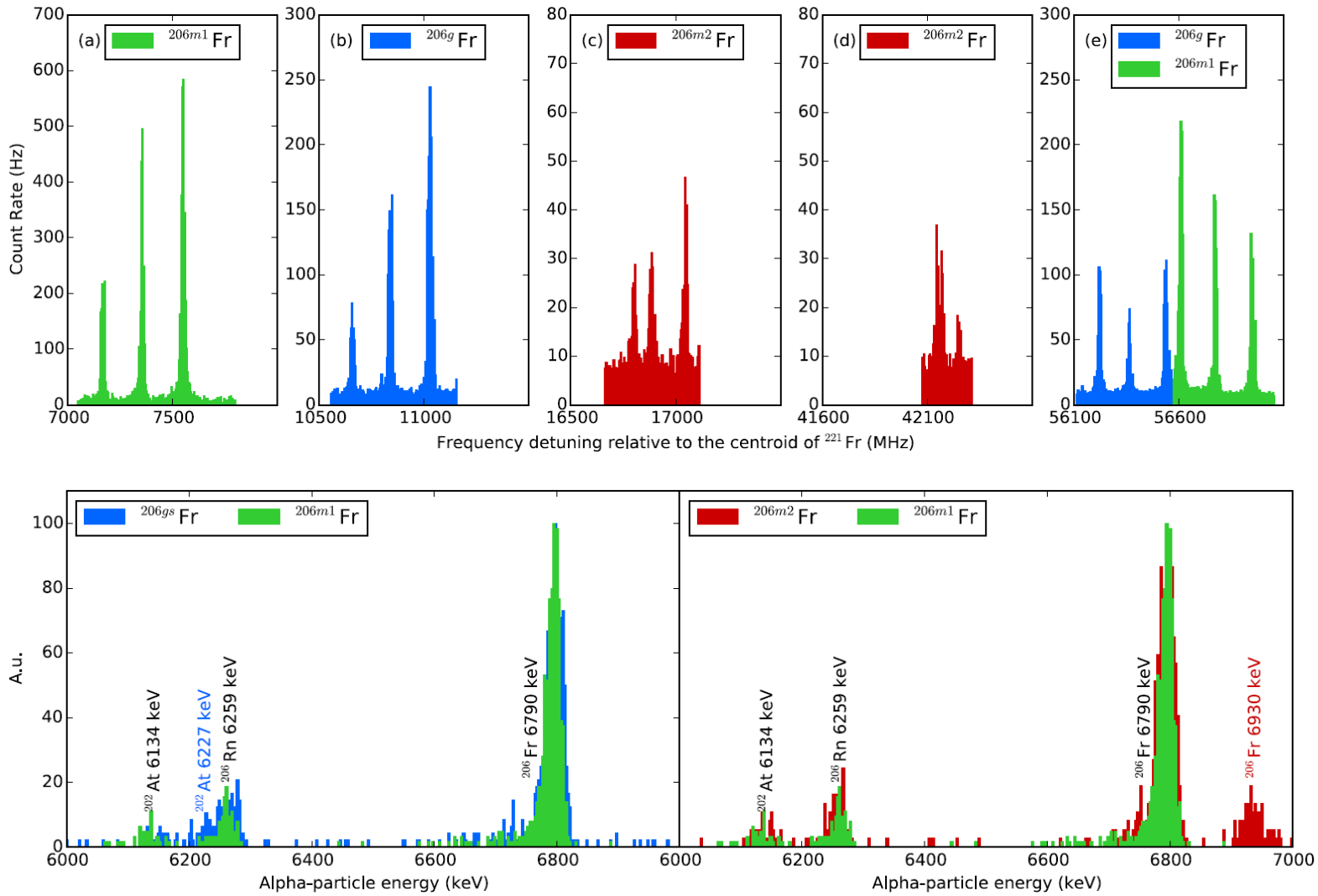
- Near-degenerate states can be de-tangled using laser-assisted nuclear decay spectroscopy.



IS471: Neutron deficient

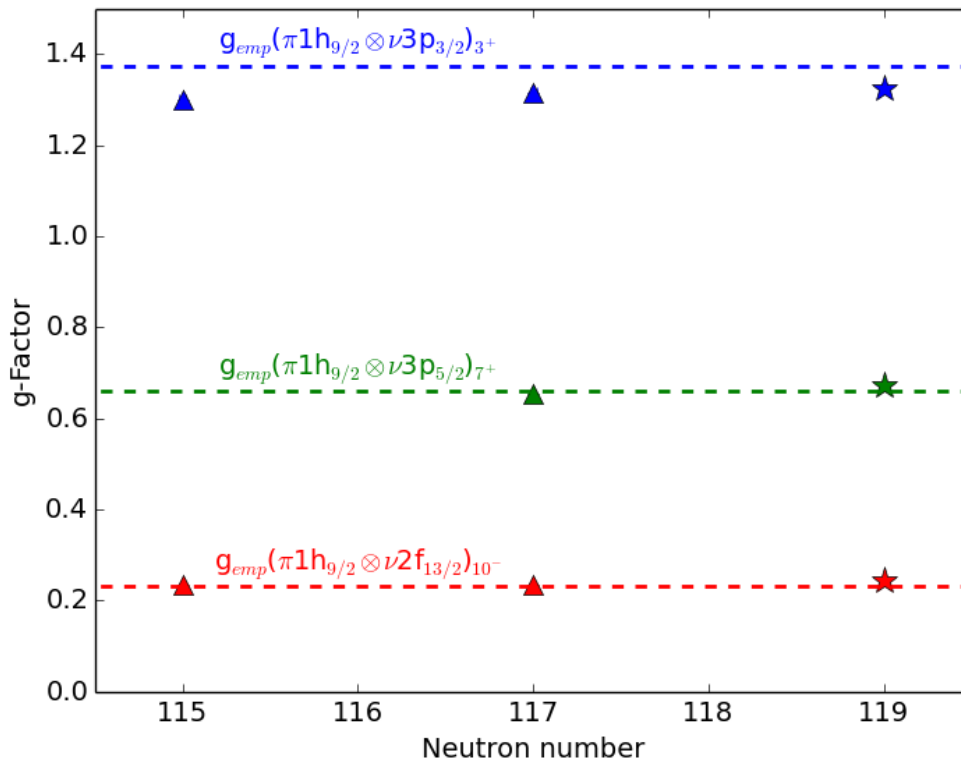


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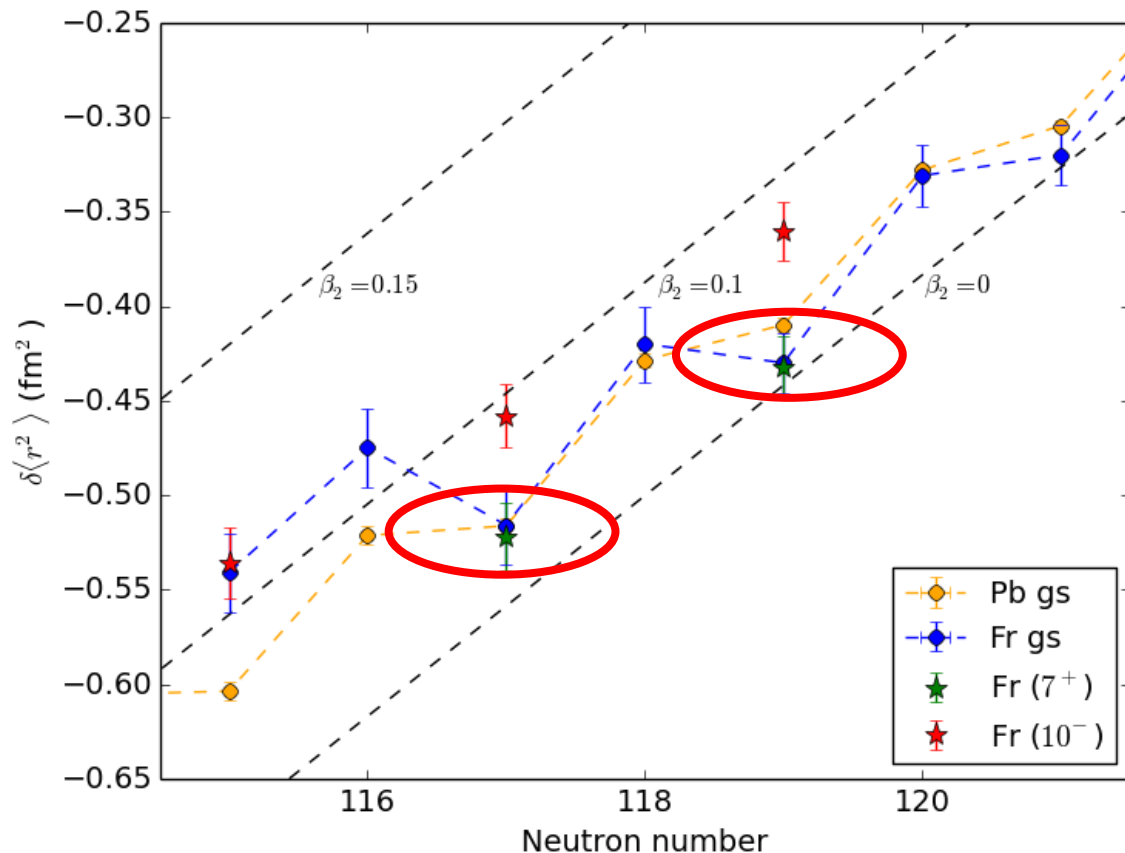
IS471: Neutron deficient

$$g_{\text{emp}} = \frac{1}{2} \left[g_p + g_n + (g_p - g_n) \frac{j_p(j_p + 1) - j_n(j_n + 1)}{I(I + 1)} \right]$$



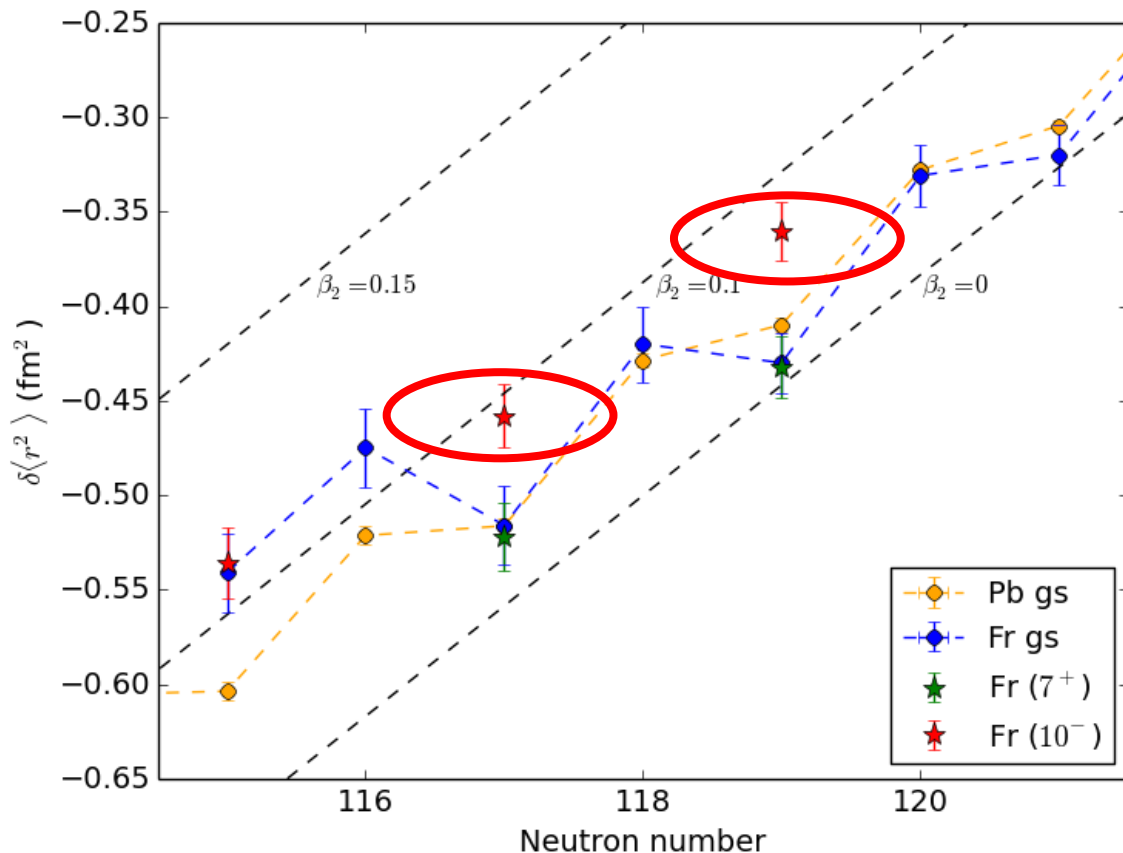
- g-factor values are consistent with the additivity rule above.
- g_{emp} values are calculated using magnetic moments of isotopes near to those being measured.

IS471: Neutron deficient



- 3⁺ and 7⁺ states exhibit spherical-like behaviour (in accordance with the Pb chain).

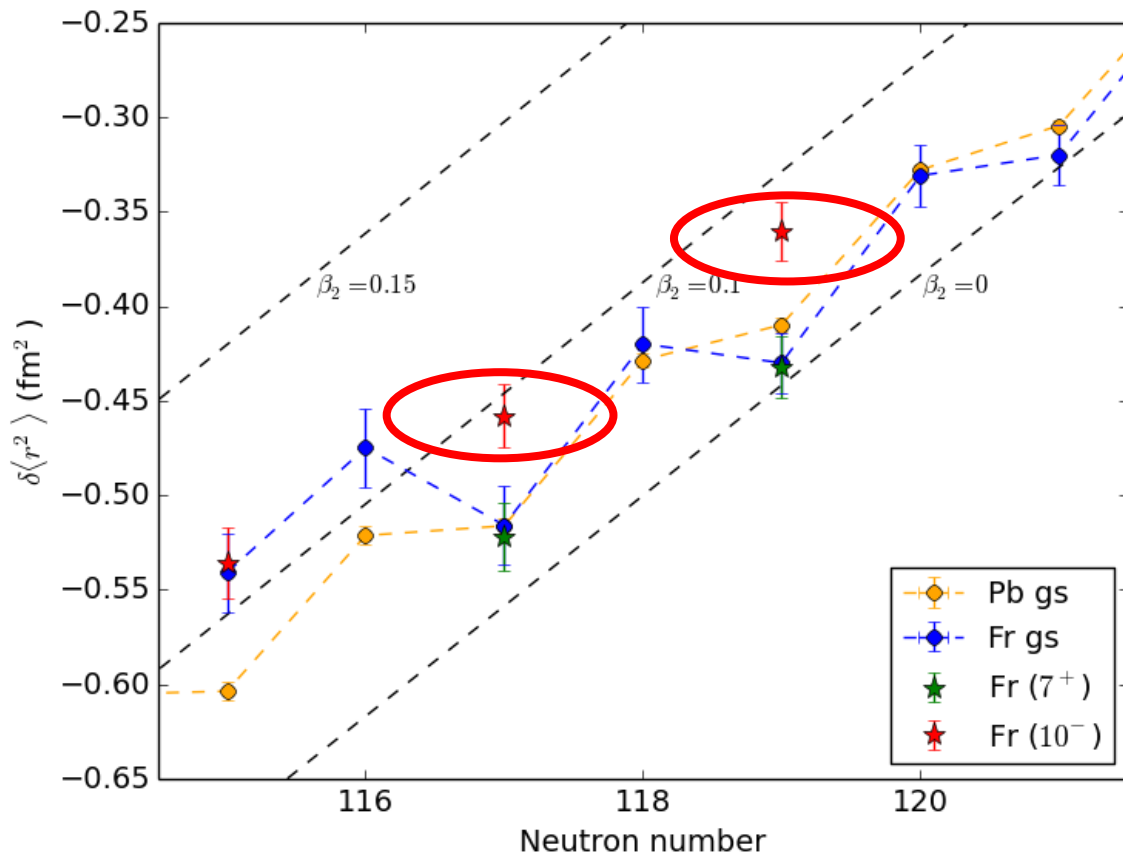
IS471: Neutron deficient



- 3^+ and 7^+ states exhibit spherical-like behaviour (in accordance with the Pb chain).
- However, the 10^- states are deformed.

IS471: Neutron deficient

Evidence of shape coexistence!

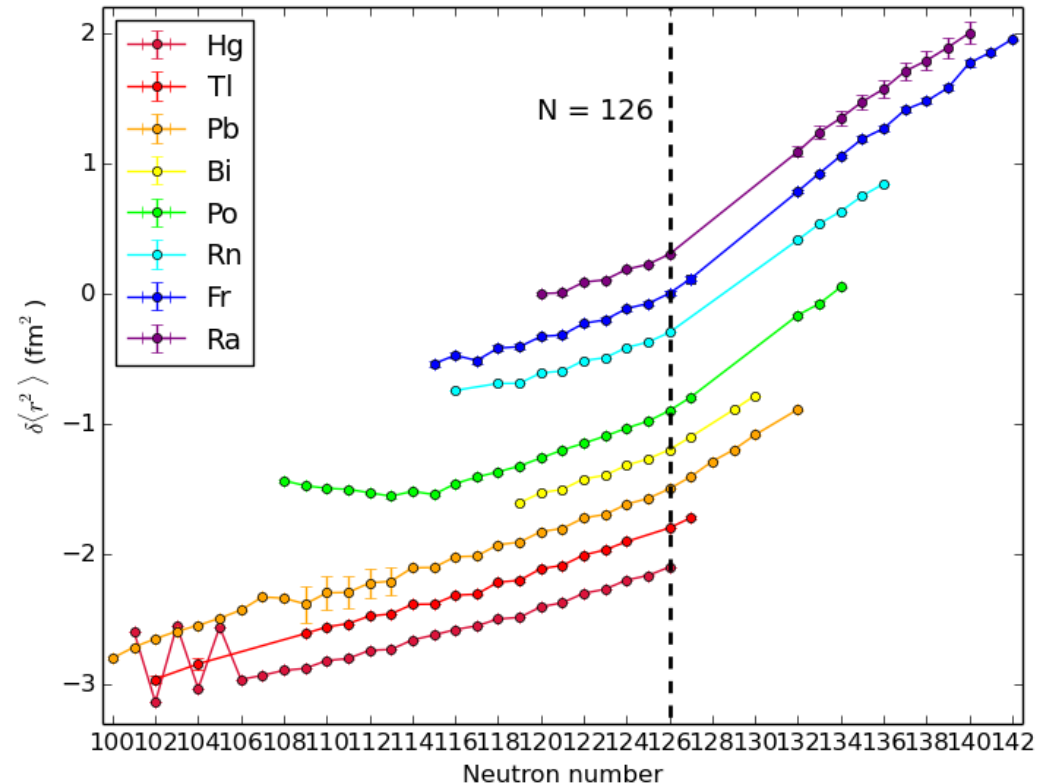


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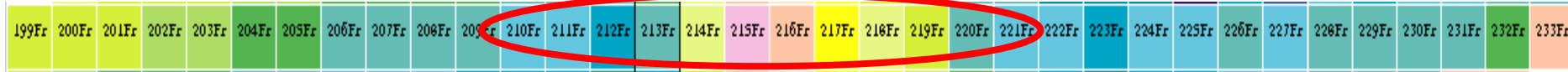
IS471: N = 126 shell closure



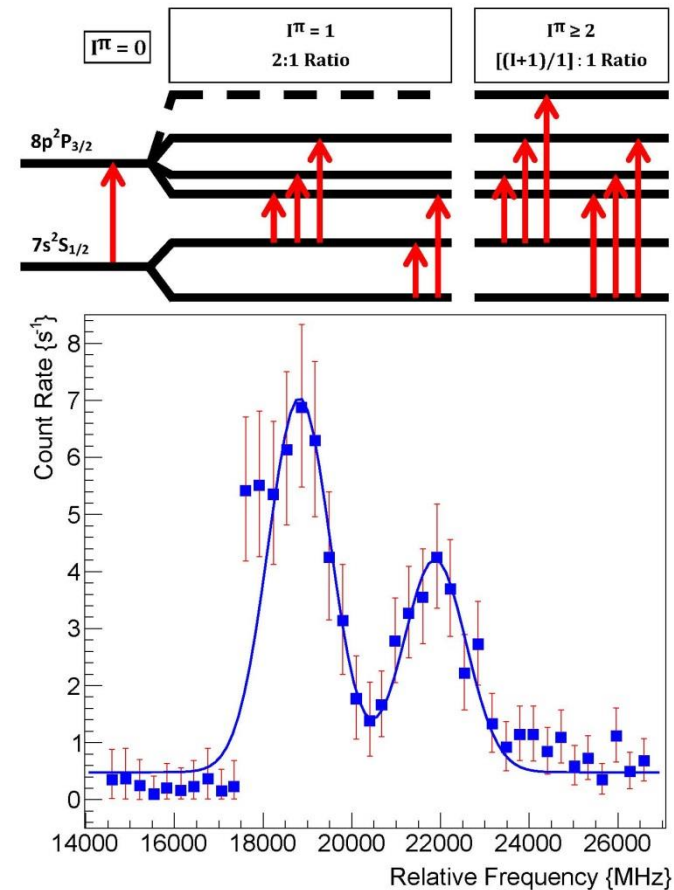
- The behaviour of these isotopes are dependent on their proximity to the N = 126 shell closure.
- In particular, the kink in the mean-square charge-radii values cannot be fully explained.



IS471: N = 126 shell closure

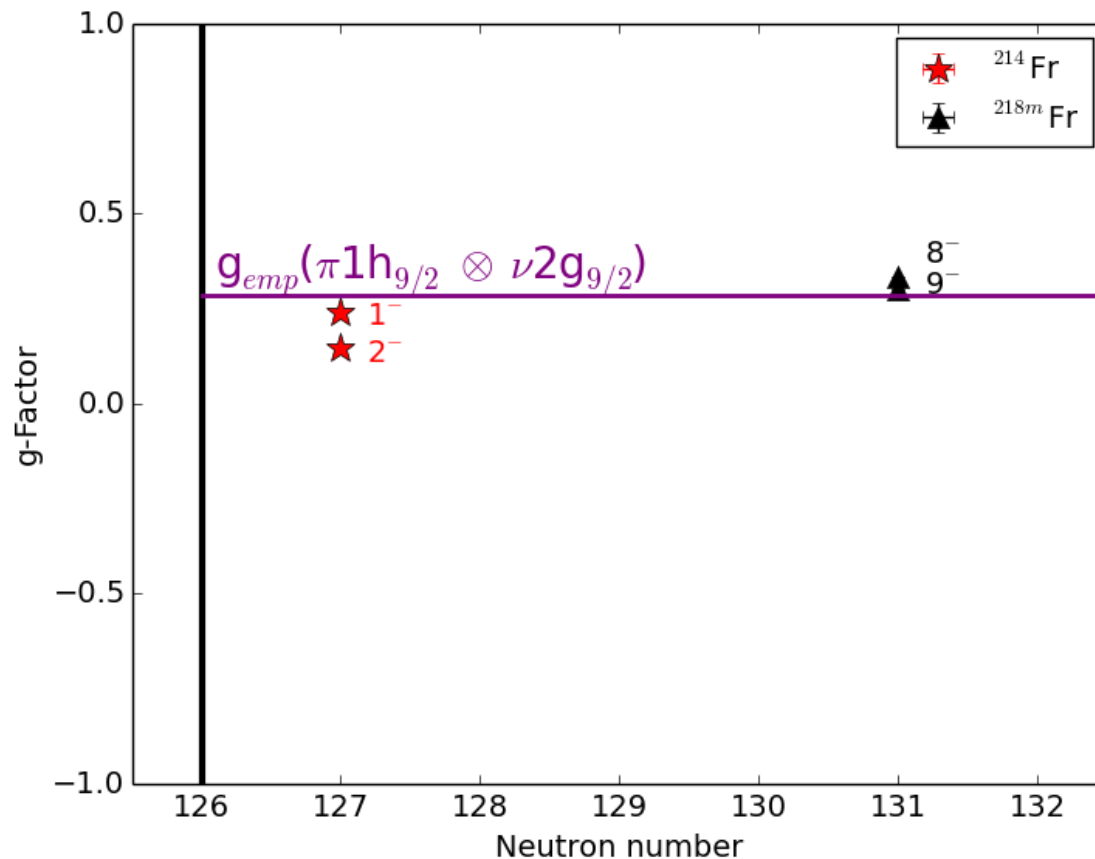


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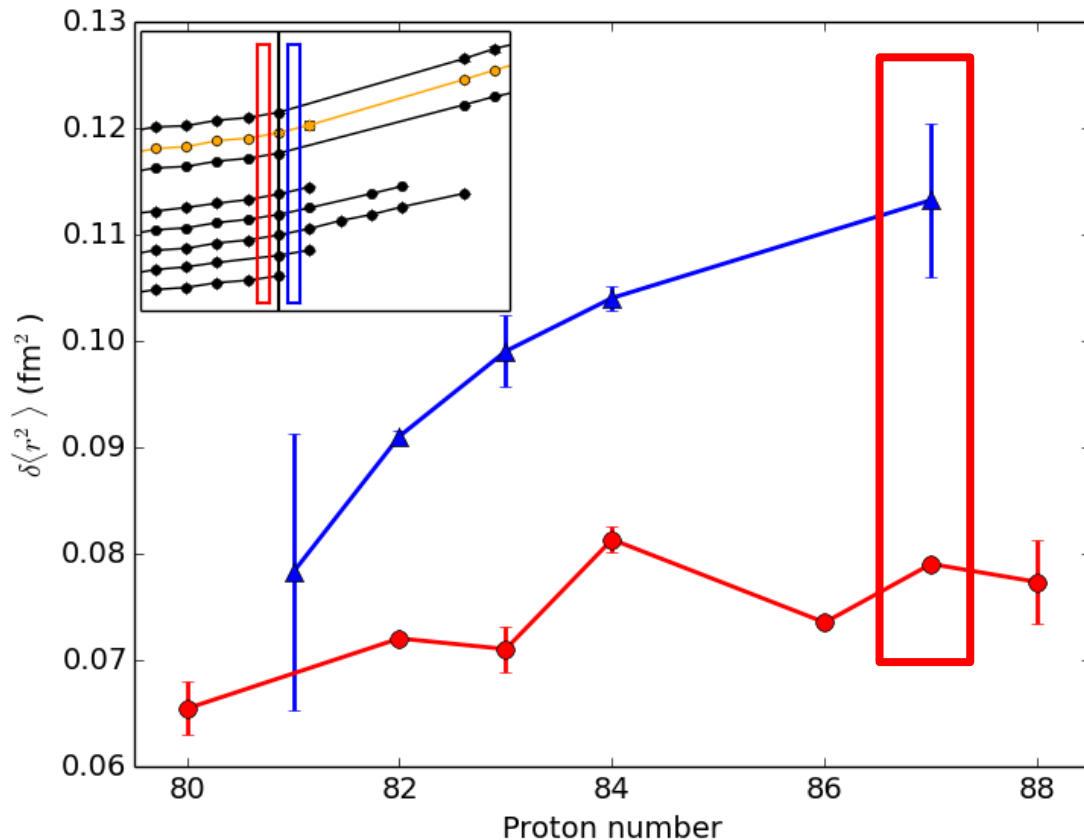
IS471: N = 126 shell closure

$$g_{\text{emp}} = \frac{1}{2} \left[g_p + g_n + (g_p - g_n) \frac{j_p(j_p + 1) - j_n(j_n + 1)}{I(I + 1)} \right]$$



- As with ^{206}Fr , g-factor values for ^{214}Fr are consistent with the additivity relation.
- Suggests a tentative spin assignment of (1^-) .

IS471: N = 126 shell closure



- As the $N = 126$ shell closure is crossed, a kink is seen in the charge-radii values, depicted by gradient increase.

199Fr	200Fr	201Fr	202Fr	203Fr	204Fr	205Fr	206Fr	207Fr	208Fr	209Fr	210Fr	211Fr	212Fr	213Fr	214Fr	215Fr	216Fr	217Fr	218Fr	219Fr	220Fr	221Fr	222Fr	223Fr	224Fr	225Fr	226Fr	227Fr	228Fr	229Fr	230Fr	231Fr	232Fr	233Fr
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IS471: $N = 126$ shell closure



PRL **110**, 032503 (2013)

PHYSICAL REVIEW LETTERS

week ending
18 JANUARY 2013

Charge Radius Isotope Shift Across the $N = 126$ Shell Gap

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(Received 9 October 2012; revised manuscript received 26 November 2012; published 15 January 2013)

We revisit the problem of the kink in the charge radius shift of neutron-rich even isotopes near the $N = 126$ shell closure. We show that the ability of a Skyrme force to reproduce the isotope shift is determined by the occupation of the neutron $1i_{11/2}$ orbital beyond $N = 126$ and the corresponding change it causes to deeply-bound protons orbitals with a principal quantum number of 1. Given the observed position of the single-particle energies, one must either ensure occupation is allowed through correlations, or not demand that the single-particle energies agree with experimental values at the mean-field level.

DOI: [10.1103/PhysRevLett.110.032503](https://doi.org/10.1103/PhysRevLett.110.032503)

PACS numbers: 21.10.Ft, 21.60.Jz, 21.30.-x, 27.80.+w

The evolution of charge radii across the isotope chart is one of the most basic nuclear structure observables, as it provides a particularly useful characterization of the proton distribution that can be accessed by a variety of experiments [1]. We define the charge radius isotope shift as the difference between the mean squared charge radius, $\langle r_{\text{ch}}^2 \rangle$, of a series of isotopes and that of a given reference isotope

Letter, we propose a new mechanism to explain the existence of the kink in lead isotope shifts using density functional calculations supplemented by pairing effects.

On the theoretical side, mean-field models, or, equivalently, density functional theories have been widely applied to the systematic study of all observed and hypothesized nuclear isotopes [8]. Both the Skyrme-Hartree-Fock and

IS471: $N = 126$ shell closure



PHYSICAL REVIEW C 91, 021302(R) (2015)

Effects of three-nucleon spin-orbit interaction on isotope shifts of Pb nuclei

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(Received 4 December 2014; published 18 February 2015)

We investigate the effects of the $3N$ interaction, which effectively adds a density-dependent term to the LS channel, on the isotope shifts of the Pb nuclei. With the strength so as to keep the ℓs splitting of the single-nucleon orbits, the density dependence in the LS channel tends to shrink the wave functions of the $j = \ell + 1/2$ orbits while making the $j = \ell - 1/2$ functions distribute more broadly. Thereby the kink in the isotope shifts of the Pb nuclei at $N = 126$ becomes stronger, owing to the attraction from neutrons occupying $0i_{11/2}$ in $N > 126$. The density dependence in the LS channel enables us to reproduce the data of the isotope shifts by the Hartree-Fock-Bogoliubov calculations in a long chain of neutron numbers, even without degeneracy between the $n1g_{9/2}$ and $n0i_{11/2}$ levels. We exemplify it by the semirealistic M3Y-P6 interaction.

Summary

- Recent work on francium isotopes at the CRIS setup have revealed a variety of nuclear features, analogous to that exhibited in other isotope chains.
- Work on $^{204,206}\text{Fr}$ revealed existence of shape coexistence in the neutron-deficient isotopes.
- A kink in the charge-radii value is seen in ^{214}Fr past the $N = 126$ shell closure, however more work needs to be done to interpret this mechanism.